

# Quantum alternation

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# Outline

- 1 Introduction
- 2 Basic background
- 3 Superoperators: Kraus, Choi and Stinespring
- 4 Classical control and quantum data
- 5 Quantum control: ideas
- 6 Quantum control: semantics
- 7 Conclusions

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- Measurement calculus: low-level, close to implementation.
- Selinger's Quantum Programming Language: Quantum data and classical control.
- There are more.



# Example

## Simple program

```
input  $b$ :bit;  
input  $p, q$ :qbit;  
 $b := \text{measure } p$ ;  
if  $b$  then  $q := N(q)$  else  $p := N(p)$ ;  
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- $N$  is the **NOT** operation on a qubit.
- **bit** and **qbit** separate datatypes.
- The conditional is based on a classical boolean.

# What about quantum alternation?

## Simple program

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- Quantum alternation is problematic in general.

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- Given a cone we can define a notion of *positive* element by saying  $x$  is positive if  $x \in C$ .
- We induce a partial order  $\leq_C$  by  $x \leq_C y$  if  $y - x \in C$ .



# Positive operators

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- A **positive map** is a map from  $\mathcal{B}(\mathcal{H})$  to itself such that it takes positive operators to positive operators.

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- Maps describing physical processes (e.g. channels) must be completely positive maps (cp maps).
- A **superoperator** is a cp map that is also trace non-increasing.

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- If  $n = m$  (square matrices) we write  $M_n$ .
- We write  $CP(M_n, M_k)$  for completely positive maps from  $M_n$  to  $M_k$ .

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- A  **$*$ -algebra** is an algebra equipped with a unary operation  $*$  such that: (i)  $a^{**} = a$ , (ii)  $(ab)^* = b^* a^*$  and (iii)  $(\lambda a)^* = \bar{\lambda} a^*$ , where  $\lambda \in \mathbb{C}$ .

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- The matrix algebras  $M_n$  are all  $C^*$ -algebras with the  $*$  being  $\dagger$  (adjoint).
- The bounded operators on a Hilbert space form a  $C^*$ -algebra.

## About $C^*$ -algebras

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- Every commutative unital  $C^*$ -algebra is isomorphic to the set of continuous functions on a compact Hausdorff space (Gelfand duality).

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- However, abstract  $C^*$ -algebras can be *represented* in a concrete way as a subalgebra of  $\mathcal{B}(\mathcal{H})$ .
- A **representation** of a  $C^*$ -algebra  $\mathcal{A}$  is a homomorphism  $\pi : \mathcal{A} \rightarrow \mathcal{B}(\mathcal{H})$  for some Hilbert space.

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- Choi: The action of  $\phi \in CP(M_n, M_k)$  can be given explicitly as a matrix in  $M_{nk}$  depending on the particular Kraus decomposition.
- Stinespring: For any completely positive map  $\theta : \mathcal{A} \rightarrow \mathcal{B}(\mathcal{H})$  there is a triple  $(\pi, V, \mathcal{K})$  where  $\mathcal{K}$  is a Hilbert space,  $\pi : \mathcal{A} \rightarrow \mathcal{B}(\mathcal{K})$  is a representation and  $V : \mathcal{H} \rightarrow \mathcal{K}$  such that

$$\theta(a) = V^\dagger \pi(a) V.$$

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- For quantum information theory this tells one that any completely positive map can be realized as a unitary on an expanded space: purification.
- If  $\theta \in CP(M_n, M_k)$  then the minimal Stinespring representation is in  $M_m$  where  $m \leq n^2k$ .

# Stinespring to Kraus

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- Let  $\mathcal{E} : \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{K})$  be a superoperator.
- By Stinespring, there exists an ancilla  $\mathcal{A}$  and an operator  $V : \mathcal{K} \rightarrow \mathcal{H} \otimes \mathcal{A}$  such that

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- Choose a basis  $\{e_i\}_{i=1}^k$  for  $\mathcal{A}$  and define  $V_i : \mathcal{K} \rightarrow \mathcal{H}$  by

$$\forall \psi \in \mathcal{K}, V\psi = \sum_{i=1}^k (V_i\psi) \otimes e_i.$$

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- Let  $\mathcal{H}$  and  $\mathcal{K}$  be two finite-dimensional Hilbert spaces and  $\mathcal{B}(\mathcal{H}), \mathcal{B}(\mathcal{K})$  the Banach algebras of bounded linear operators.
- Let  $\mathcal{E} : \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{K})$  be a superoperator.
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- The  $V_i$  give a Kraus representation for  $\mathcal{E}$ .



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- A function from a dcpo to another dcpo is called **Scott continuous** if it preserves lubs of increasing sequences.



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- Recursion can implement iteration but not the other way around.

# What do we want?

- Suppose we have a qubit  $q$  and two superoperators  $S, T : \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{K})$  then the quantum alternation  $(qAlt)(q; S, T)$  should be a superoperator from  $\mathcal{B}(\mathcal{Q} \otimes \mathcal{H}) \rightarrow \mathcal{B}(\mathcal{Q} \otimes \mathcal{K})$ .

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- We want the operation to be monotone so we can use this inside recursions.



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- Action:  $(\sum_i e_i \otimes \psi_i) \mapsto (\sum_i e_i \otimes U_i \psi_i)$ .

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- Very useful for describing algorithms especially if there are only unitaries.

## Examples II: Deutsch's algorithm

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**new qbit**  $x, y$

$$x^* = H$$

$$y^* = N; H$$

**if**  $x$  **then**  $y^* = U_0$  **else**  $y^* = U_1$

$$x^* = H$$

# Example III: Quantum Fourier transform



**for**  $i = 1$  **to**  $n$  **do**

$q_i$   $\ast = H$

**for**  $k = 2$  **to**  $n - i + 1$  **do**

**if**  $q_{k+i-1}$  **then skip else**  $q_i \ast = R_k$

Here  $R_k$  is the phase shift gate defined by  $R_k = \Pi_0 + e^{i\theta}\Pi_1$  with  $\theta = 2\pi/2^k$ .

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- can we extend it to quantum operations that are not unitaries?



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- A particular Kraus form comes from a particular choice of basis of the environment, as we saw.
- A basis corresponds to a particular choice of measurement. Thus the particular Kraus representation is dictated by how the experimenter chooses to describe the environment.

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- We give compositional semantics but in terms of specific choices of Kraus operators, we do not try to give compositional superoperator semantics.

# Quantum alternation of unitaries

Given unitary operators  $U, V$  on  $\mathcal{H}$  and a qubit  $q$  (space  $\mathcal{Q}$ ) we define

$$|0\rangle\langle 0| \otimes U + |1\rangle\langle 1| \otimes V = \begin{pmatrix} U & 0 \\ 0 & V \end{pmatrix}$$

as the quantum alternation of  $U$  and  $V$ .

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- This defines a superoperator

$$\mathcal{S}(\rho) = \sum_{i,j} K_{i,j}^* \rho K_{i,j}.$$

## What Stinespring says

If one looks at the Stinespring dilation corresponding to the above construction we see that the ancilla spaces (environments) of the two Kraus forms are tensored together.

# Kraus semantics

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- Applying a unitary

$$\llbracket q^* = U \rrbracket = \{U\}.$$



# More semantics

- Measure  $q$ , this has type  $\tau \rightarrow \tau \oplus \tau$

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- We do not give semantics for loops and conditionals.

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- This example arose from discussions with Mingsheng Ying and Yuan Feng at UTS Sydney based on an example due to Nengkun Yu.
- One can think of quantum alternation as an algorithmic notation, it is not clear what it means *physically*.



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- By explicit calculation we can show that  $R' \not\leq R$ .



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- The BARN then gives a Kraus decomposition.
- One can give a denotational semantics based on these “canonical” Kraus forms but there is little reason to think that this has physical significance.

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- They used a notion of “orthogonality” and only allow orthogonal terms to be put in quantum alternation.
- However, they did not give complete rules. For example, one cannot nest quantum conditionals.

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- Defined a superoperator semantics and noted lack of compositionality.
- Implicit in their superoperator semantics is our Kraus semantics.
- Perhaps one should view the superoperator semantics as an *abstract interpretation* of the Kraus semantics.
- Did not note non-monotonicity but had a different approach to recursion based on Fock space [Ying 2015].

# Summary

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# Summary

- Quantum alternation is troublesome: non-compositional and non-monotone.
- Is it a sensible thing to even consider? It came from programming languages without thinking about physics.
- One should look at real physical situations, e.g. Mach-Zehnder interferometers and extract a notion of quantum alternation. Hines-Scott develop a notion of conditional iteration along these lines.
- Perhaps quantum alternation and recursion is not allowed in nature!

Thank you!